

Modelling Electroweak Interactions at High Energies

Wolfgang Kilian

University of Siegen

Peking University

May 2015

Higgs and Vector-Boson Scattering

$$O(E^4) \text{ (diagram)} + O(E^4) \text{ (diagram)} + O(E^2) \text{ (diagram)} = O(1)$$

Higgs exchange cancels the E^2 rise exactly (in the SM):
 the **Minimal SM Higgs Sector**.

Higgs and Vector-Boson Scattering

$$O(E^4) \text{ (diagram)} + O(E^4) \text{ (diagram)} + O(E^2) \text{ (diagram)} = O(1)$$

Higgs exchange cancels the E^2 rise exactly (in the SM):
the Minimal SM Higgs Sector.

Discoveries

1. Higgs production in WW fusion: the Higgs **boson** exists.
2. SM confirmed in VBS: the Higgs **mechanism** works as expected.

Future Expectation for VBS

If SM is true,

Future Expectation for VBS

If SM is true, VBS **amplitude is bounded** and small: m_H^2/v^2 .

LHC:

Production cross section falls off with increasing effective energy, i.e., invariant mass of the WW pair system.

NLO: some logarithmic corrections.

Future Expectation for VBS

If SM is true, VBS amplitude is bounded and small: m_H^2/v^2 .

LHC:

Production cross section falls off with increasing effective energy, i.e., invariant mass of the WW pair system.

NLO: some logarithmic corrections.

No problem with **unitarity**, of course.

And What If Not?

Two classes of modifications to the SM (or mixture):

1. New weakly interacting particles, **direct production**. Example: 2HDM

And What If Not?

Two classes of modifications to the SM (or mixture):

1. New weakly interacting particles, direct production. Example: 2HDM
2. **Small deviations** from the SM prediction (**linear Higgs rep.**)

And What If Not?

Two classes of modifications to the SM (or mixture):

1. New weakly interacting particles, direct production. Example: 2HDM
2. Small deviations from the SM prediction (linear Higgs rep.)

⇒ large effect at multi-TeV energy (SPPC, CLIC)

And What If Not?

Two classes of modifications to the SM (or mixture):

1. New weakly interacting particles, direct production. Example: 2HDM
2. Small deviations from the SM prediction (linear Higgs rep.)

⇒ large effect at multi-TeV energy (SPPC, CLIC)

Effective Field Theory

- ▶ Add higher-dimensional operators to the SM Lagrangian.
- ▶ Use only SM fields, respect SM gauge invariance
- ▶ Operator of dimension n carries prefactor $1/\Lambda^{n-4}$

And What If Not?

Two classes of modifications to the SM (or mixture):

1. New weakly interacting particles, direct production. Example: 2HDM
2. Small deviations from the SM prediction (linear Higgs rep.)

⇒ large effect at multi-TeV energy (SPPC, CLIC)

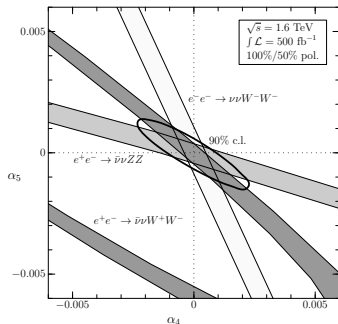
Effective Field Theory

- ▶ Add higher-dimensional operators to the SM Lagrangian.
- ▶ Use only SM fields, respect SM gauge invariance
- ▶ Operator of dimension n carries prefactor $1/\Lambda^{n-4}$

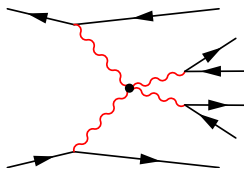
$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \sum_{d=5}^{\infty} \frac{1}{\Lambda^{n-4}} \mathcal{O}_n$$

Vector-Boson Scattering at high energies,

E. Boos, H.-J. He, WK, A. Pukhov, C.-P. Yuan, P.M. Zerwas, PRD 57 (1998)



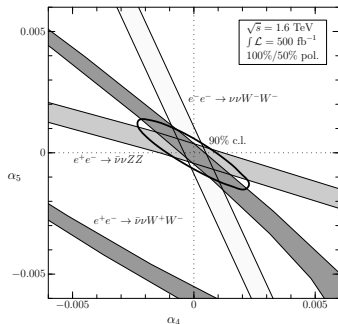
$$e^-e^+ \rightarrow \nu\bar{\nu}W^+W^- \rightarrow \nu\bar{\nu} + 4f$$



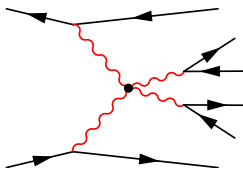
Closer study required a new Monte-Carlo Simulation program: WHIZARD

Vector-Boson Scattering at high energies,

E. Boos, H.-J. He, WK, A. Pukhov, C.-P. Yuan, P.M. Zerwas, PRD 57 (1998)

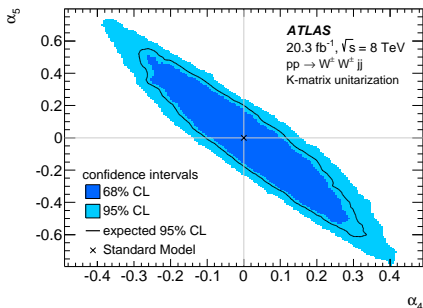
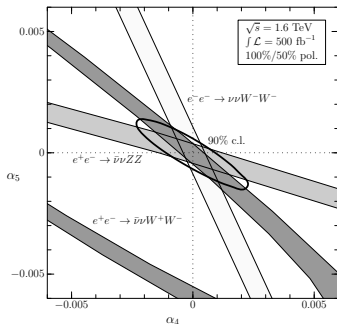


$$e^-e^+ \rightarrow \nu\bar{\nu}W^+W^- \rightarrow \nu\bar{\nu} + 4f$$



Closer study required a new Monte-Carlo Simulation program: WHIZARD Vector-Boson Scattering at high energies,

E. Boos, H.-J. He, WK, A. Pukhov, C.-P. Yuan, P.M. Zerwas, PRD 57 (1998)

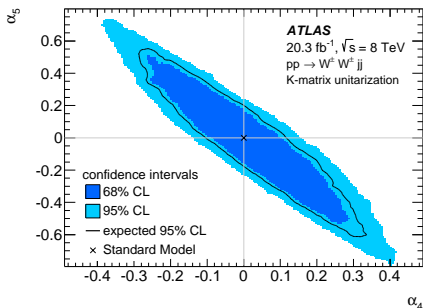
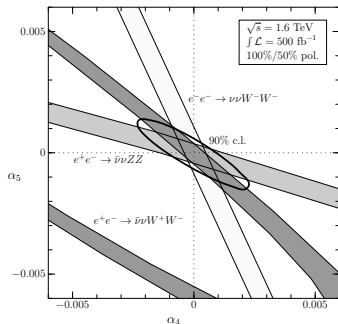


Discovery of Vector-Boson Scattering

ATLAS collaboration, PRL 113 (2014)

Closer study required a new Monte-Carlo Simulation program: WHIZARD Vector-Boson Scattering at high energies,

E. Boos, H.-J. He, WK, A. Pukhov, C.-P. Yuan, P.M. Zerwas, PRD 57 (1998)



Discovery of Vector-Boson Scattering

ATLAS collaboration, PRL 113 (2014) — using WHIZARD

Tool for Calculation and Simulation

WHIZARD is a universal Monte Carlo for high-energy processes
(first version 1999, currently 2.2.6)

- ▶ tailored for lepton colliders (LEP, ILC, CEPC, CLIC)
- ▶ applied also for hadron colliders (ATLAS/CMS, CPPC)

The WHIZARD Team

- ▶ University of Siegen: Wolfgang Kilian
- ▶ DESY (Hamburg): Jürgen Reuter
- ▶ University of Würzburg: Thorsten Ohl

New Study of WW Scattering

WK, T. Ohi, J. Reuter, M. Sekulla, arXiv:1408.6207 (PRD, to appear)

- ▶ Investigate quartic anomalous couplings of vector bosons
- ▶ Match to EFT with Higgs (linear representation)
- ▶ Extrapolate to high energies without violating unitarity in the calculation
- ▶ Implement as non-Lagrangian model in WHIZARD
- ▶ Evaluate for full SM particle set, observables LHC processes

Concrete Examples:

Anomalous Interactions

$$\mathcal{L}_{HD} = F_{HD} \operatorname{tr} \left[\mathbf{H}^\dagger \mathbf{H} - \frac{v^2}{4} \right] \cdot \operatorname{tr} \left[(\mathbf{D}_\mu \mathbf{H})^\dagger (\mathbf{D}^\mu \mathbf{H}) \right] \quad HVV \quad D = 6$$

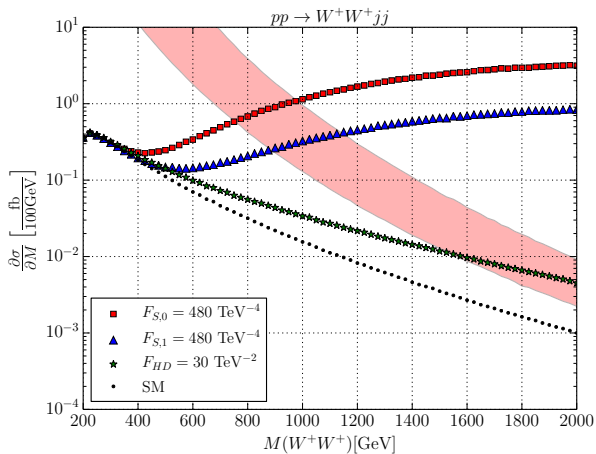
$$\mathcal{L}_{S,0} = F_{S,0} \operatorname{tr} \left[(\mathbf{D}_\mu \mathbf{H})^\dagger \mathbf{D}_\nu \mathbf{H} \right] \cdot \operatorname{tr} \left[(\mathbf{D}^\mu \mathbf{H})^\dagger \mathbf{D}^\nu \mathbf{H} \right] \quad VVV \quad D = 8$$

$$\mathcal{L}_{S,1} = F_{S,1} \operatorname{tr} \left[(\mathbf{D}_\mu \mathbf{H})^\dagger \mathbf{D}^\mu \mathbf{H} \right] \cdot \operatorname{tr} \left[(\mathbf{D}_\nu \mathbf{H})^\dagger \mathbf{D}^\nu \mathbf{H} \right] \quad VVV \quad D = 8$$

Linear Higgs/Goldstone Field Representation:

$$\mathbf{H} = \frac{1}{2} \begin{pmatrix} v + h - iw^3 & -i\sqrt{2}w^+ \\ -i\sqrt{2}w^- & v + h + iw^3 \end{pmatrix} \cdot \quad (1)$$

Nice, but...



Calculation: WHIZARD

What happened?

Gauge invariance + Higgs exchange
remove **two** orders of the Taylor expansion.

⇒ Effect of anomalous couplings rapidly rises with energy. ($D = 8$ operators!) **cancels the PDF suppression**

What happened?

Gauge invariance + Higgs exchange
remove two orders of the Taylor expansion.

- ⇒ Effect of anomalous couplings rapidly rises with energy. ($D = 8$ operators!) cancels the PDF suppression
- ⇒ Window (in energy) where effective theory is useful for describing deviations at the LHC: absent.

Basically, **forget about (perturbative) quantum field theory?**

This is not the same situation as in VB pair production.

What happened?

Gauge invariance + Higgs exchange
remove two orders of the Taylor expansion.

- ⇒ Effect of anomalous couplings rapidly rises with energy. ($D = 8$ operators!) cancels the PDF suppression
- ⇒ Window (in energy) where effective theory is useful for describing deviations at the LHC: absent.

Basically, forget about (perturbative) quantum field theory?

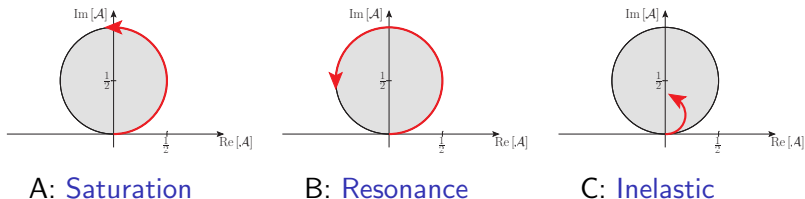
This is not the same situation as in VB pair production.

[There are **perturbative** models, e.g, the 2HDM. But they access only a **small fraction** of the conceivable Model Space.]

Unitarity

The scattering of w, z is a (quasi-) elastic process. Properly diagonalized (isospin I , spin J) and normalized, the partial-wave amplitudes must lie on the Argand Circle.

Possibilities



Unitarization

Unitarization

There are zillions of papers that investigate this problem.

- ▶ Heavy Higgs as Unitarization
- ▶ K-Matrix Unitarization
- ▶ Padé Unitarization
- ▶ Inverse Amplitude Method
- ▶ $O(N)$ Model Unitarization
- ▶ N/D Method
- ▶ ...

Unitarization

There are zillions of papers that investigate this problem.

- ▶ Heavy Higgs as Unitarization
- ▶ K-Matrix Unitarization
- ▶ Padé Unitarization
- ▶ Inverse Amplitude Method
- ▶ $O(N)$ Model Unitarization
- ▶ N/D Method
- ▶ ...

Small caveat: 99 % of those papers **don't have a light Higgs.**

Unitarization

There are zillions of papers that investigate this problem.

- ▶ Heavy Higgs as Unitarization
- ▶ K-Matrix Unitarization
- ▶ Padé Unitarization
- ▶ Inverse Amplitude Method
- ▶ $O(N)$ Model Unitarization
- ▶ N/D Method
- ▶ ...

Small caveat: 99 % of those papers don't have a light Higgs.

Which makes a difference.

Unitarization after 2012

Repeat the game with light Higgs?

Unitarization methods are tailored for the quasi-elastic WW system, not for arbitrary processes.

Unitarization after 2012

Repeat the game with light Higgs?

Unitarization methods are tailored for the quasi-elastic WW system, not for arbitrary processes.

- ▶ **measure** low-energy parameters
- ▶ **extrapolate**, using analytic properties and assumptions
- ▶ get a **prediction**.

Unitarization after 2012

Repeat the game with light Higgs?

Unitarization methods are tailored for the quasi-elastic WW system, not for arbitrary processes.

- ▶ measure low-energy parameters
- ▶ extrapolate, using analytic properties and assumptions
- ▶ get a prediction.

Do we want a prediction with assumptions?

Unitarization after 2012

Repeat the game with light Higgs?

Unitarization methods are tailored for the quasi-elastic WW system, not for arbitrary processes.

- ▶ measure low-energy parameters
- ▶ extrapolate, using analytic properties and assumptions
- ▶ get a prediction.

Do we want a prediction with assumptions? We want a **framework**.

K Matrix

(Heitler 1941, for QED): Cayley Transform

$$S = \frac{\mathbb{1} + iK/2}{\mathbb{1} - iK/2}, \quad \text{where } K = K^\dagger \quad \text{and} \quad S = \mathbb{1} + iT$$

The K Matrix, exactly:

$$K = \frac{T}{\mathbb{1} + iT/2}.$$

K Matrix

(Heitler 1941, for QED): Cayley Transform

$$S = \frac{\mathbb{1} + iK/2}{\mathbb{1} - iK/2}, \quad \text{where } K = K^\dagger \quad \text{and} \quad S = \mathbb{1} + iT$$

The K Matrix, exactly:

$$K = \frac{T}{\mathbb{1} + iT/2}.$$

The K Matrix, in Perturbation Theory:

$$K = T - \frac{i}{2}T^2 \pm \dots$$

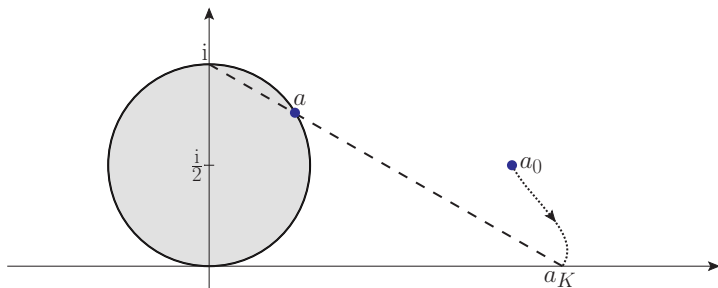
Original K Matrix algorithm (**Gupta**, for QCD/EW):

- ▶ Compute T matrix perturbatively
- ▶ Reconstruct K matrix order by order
- ▶ Insert into S matrix formula, without expanding again

This is elegant, but relies on perturbation theory.

Graphical Visualization: K Matrix

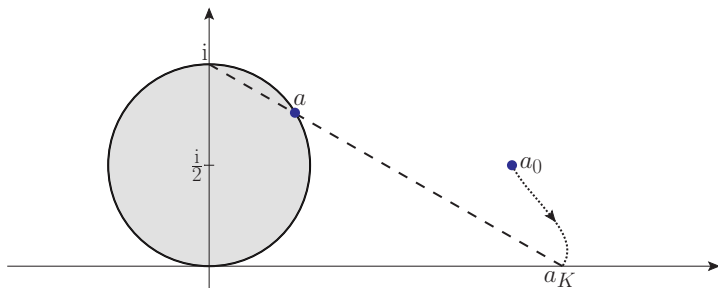
Start from **arbitrary** amplitude a_0 in perturbative expansion:



First reconstruct a_K , then compute a

Graphical Visualization: K Matrix

Start from **arbitrary** amplitude a_0 in perturbative expansion:

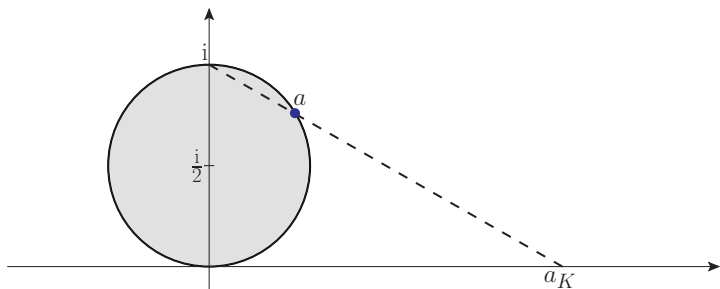


First reconstruct a_K , then compute a

Our suggestion: compute unitarized T matrix directly, without detour

Graphical Visualization: Direct T Matrix Unitarization

Start from **real** amplitude $a_0 = a_K$: Inverse stereographic projection

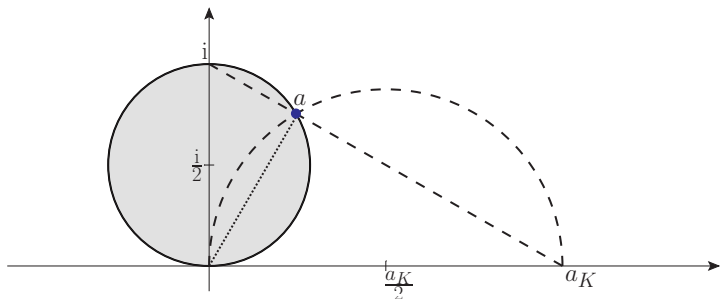


⇒ No reference to perturbative expansion

⇒ Unitary amplitude a_0 left invariant

Graphical Visualization: Direct T Matrix Unitarization

Start from **real** amplitude $a_0 = a_K$: Thales circle projection

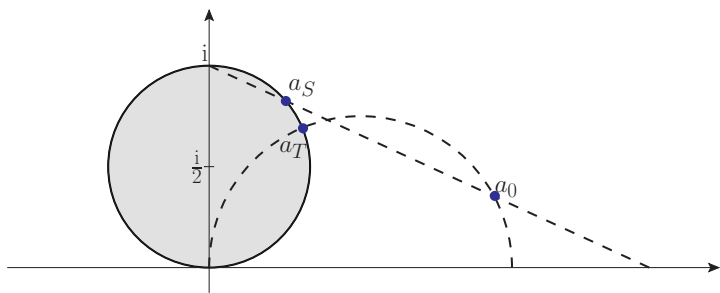


⇒ No reference to perturbative expansion

⇒ Unitary amplitude a_0 left invariant

Graphical Visualization: Direct T Matrix Unitarization

Start from **complex** amplitude a_0 :



- ⇒ No reference to perturbative expansion
- ⇒ Unitary amplitude a_0 left invariant
- ⇒ But **scheme dependence** for complex a_0

Linear Construction “Stereographic”

$$T = \frac{\operatorname{Re} T_0}{\mathbb{1} - \frac{i}{2} T_0^\dagger}.$$

for normal matrices ($T^\dagger T = T T^\dagger$), otherwise need operator ordering

- ▶ well behaved near $T = 0$
- ▶ weird behavior for eigenvalues above $T = i$

Circular Construction “Thales”

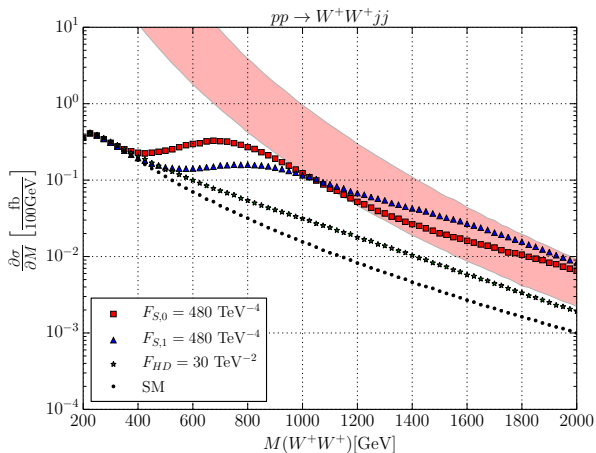
$$T = \frac{1}{\operatorname{Re} \left(\frac{1}{T_0} \right) - \frac{i}{2} \mathbb{1}} .$$

- ▶ singular at $T = 0$ (but harmless)
- ▶ well behaved above $T = i$

Algorithm

1. Start with input **model**
2. Extract strong-interaction part in Goldstone limit
3. **Unitarize** via T Matrix projection
4. Re-insert correction as form factor into Feynman rules
5. Extrapolate off-shell
6. Use in **Monte Carlo** simulation

Result: Unitarized Cross Section



Calculation: WHIZARD

And Beyond?

- ▶ Padé & Co. yield **predictions**: resonances
- ▶ work in QCD (vector dominance) ...?
- ▶ restricted to quasi-elastic scattering?

⇒ Add any additional information in T Matrix framework

Resonances and Anomalous Couplings

A resonance is a **pole** in the elastic scattering matrix:

$$A(s) = \frac{g^2}{s - \hat{m}^2} + \hat{A}_{\text{nonres}}(s)$$

The parameters g^2 and \hat{m}^2 are well defined: pole location and residue.

Resonances and Anomalous Couplings

A resonance is a **pole** in the elastic scattering matrix:

$$A(s) = \frac{g^2}{s - \hat{m}^2} + \hat{A}_{\text{nonres}}(s)$$

The parameters g^2 and \hat{m}^2 are well defined: pole location and residue.
Applying T-matrix unitarization, we get a **Breit-Wigner resonance**

$$A(s) = \frac{g^2}{s - m^2 + im\Gamma} + A_{\text{nonres}}(s)$$

Resonances and Anomalous Couplings

A resonance is a **pole** in the elastic scattering matrix:

$$A(s) = \frac{g^2}{s - \hat{m}^2} + \hat{A}_{\text{nonres}}(s)$$

The parameters g^2 and \hat{m}^2 are well defined: pole location and residue. Applying T-matrix unitarization, we get a **Breit-Wigner resonance**

$$A(s) = \frac{g^2}{s - m^2 + im\Gamma} + A_{\text{nonres}}(s)$$

At low energy, the resonant amplitude has a Taylor expansion

$$A(s) = -\frac{g^2}{m^2} + \frac{g^2}{m^4} s + \dots$$

The second term corresponds to an anomalous coupling (**matching**).

Guideline for Simplified Models

- ▶ The rise of an amplitude (anomalous coupling) may be the Taylor expansion of a resonance.
- ▶ We have no idea which resonances exist and where they come from.
- ▶ Including a resonance in the model, there still may be further sources for anomalous couplings (further resonances, $A_{\text{nonres}}(s)$, deviation from the Breit-Wigner shape, etc.)
- ▶ Beyond the resonance, the amplitude may eventually rise and need unitarization again.

Consequence:

We allow for resonances in all accessible spin/isospin channels.

We also include extra anomalous couplings.

Simplified Models: Generic Resonances

I	0	1	2
$J=0$	σ^0	.	$\phi^{--}, \phi^-, \phi^0, \phi^+, \phi^{++}$
1	.	ρ^-, ρ^0, ρ^+	.
2	f^0	.	$t^{--}, t^-, t^0, t^+, t^{++}$
...

- ▶ $I = 0$: resonant in W^+W^- and ZZ scattering
- ▶ $I = 1$: resonant in W^+Z and W^-Z scattering
- ▶ $I = 2$: resonant in W^+W^+ and W^-W^- scattering

Model Parameters

VBS, total (isospin preserved, CP, higher spin ignored):

- ▶ 5 resonances with 3 parameters each (M, g_L, g_T)
- ▶ quartic anomalous couplings of longitudinal VB
- ▶ quartic anomalous couplings of transversal VB
- ▶ quartic anomalous couplings mixing T and L

Project

WK, T. Ohl, J. Reuter, M. Sekulla, 2015 (work in progress)

- ▶ Supplement SM by generic resonances in the TeV range
- ▶ Match to EFT with Higgs (linear representation)
- ▶ Extrapolate to high energies without violating unitarity in the calculation
- ▶ Implement as non-Lagrangian model in WHIZARD
- ▶ Evaluate for full SM particle set, observables LHC processes

Breakdown of Possibilities

Breakdown of Possibilities:

1. Vector resonance (ρ)

- ▶ Parity odd, isospin 1 enhanced
- ▶ Can mix with W, Z
- ▶ Can thus couple to light-fermion currents

2. Scalar and/or tensor resonance ($\sigma/\phi/f/t$)

- ▶ C,P even (or odd)
- ▶ Isospin 0 and/or 2
- ▶ Scalar case may be weakly coupled (renormalizable)
- ▶ May mix with Higgs
- ▶ Expect scalar/tensor couplings to heavy-fermion currents (if any)

Breakdown of Possibilities

Breakdown of Possibilities:

1. Vector resonance (ρ)

- ▶ Parity odd, isospin 1 enhanced
- ▶ Can mix with W, Z
- ▶ Can thus couple to light-fermion currents

2. Scalar and/or tensor resonance ($\sigma/\phi/f/t$)

- ▶ C,P even (or odd)
- ▶ Isospin 0 and/or 2
- ▶ Scalar case may be weakly coupled (renormalizable)
- ▶ May mix with Higgs
- ▶ Expect scalar/tensor couplings to heavy-fermion currents (if any)

Let's consider the second case.

Breakdown of Possibilities

Breakdown of Possibilities:

1. Vector resonance (ρ)

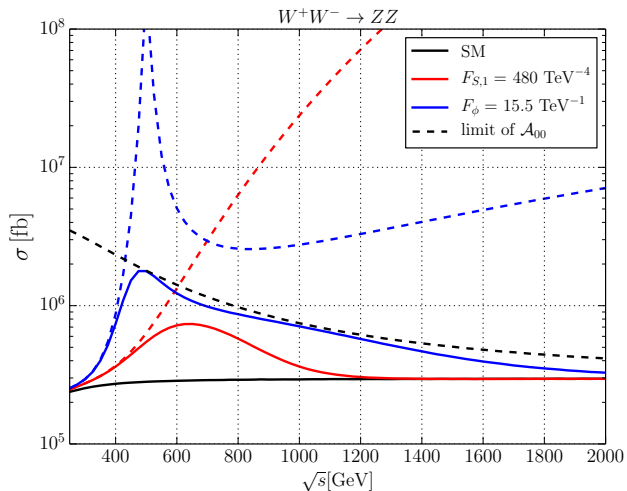
- ▶ Parity odd, isospin 1 enhanced
- ▶ Can mix with W, Z
- ▶ Can thus couple to light-fermion currents

2. Scalar and/or tensor resonance ($\sigma/\phi/f/t$)

- ▶ C,P even (or odd)
- ▶ Isospin 0 and/or 2
- ▶ Scalar case may be weakly coupled (renormalizable)
- ▶ May mix with Higgs
- ▶ Expect scalar/tensor couplings to heavy-fermion currents (if any)

Let's consider the second case. cf. extra dimensions

Example: Scalar Resonance



The Tensor Is Special

Massive spin-2 field in relativistic quantum field theory

On-shell: $SO(3)$ rotation group in rest frame: 5 polarization components

Off-shell: $SO(3) \times SO(3)$ Lorentz group: 10 components of symmetric tensor

\Rightarrow 5 components: genuine tensor, couples to current $J_{\mu\nu}$

\Rightarrow 3 components: vector, couples to divergence $\partial^\mu J_{\mu\nu}$

\Rightarrow 1 + 1 components: 2 scalars, couple to J_μ^μ and $\partial^\mu \partial^\nu J_{\mu\nu}$

Embedded in Standard Model: $\partial_\mu \rightarrow D_\mu$, and divergences don't vanish due to EWSB

The Tensor Is Special

Consequences for scattering amplitudes:

Beyond the resonance, amplitudes rise again

- ▶ Goldstones coupled to genuine tensor \Rightarrow analog to scalar, can use T matrix method for the amplitude
- ▶ Goldstones coupled to scalar \Rightarrow additional terms in T matrix
- ▶ EWSB contributions for scalar parts \Rightarrow extra terms proportional to m_W^2/M_T^2 and m_h^2/M_T^2 , non-universal
- ▶ EWSB contribution for vector parts \Rightarrow extra terms proportional to m_W^2/M_T^2 in transversal VB interactions

Have to estimate contributions that are not covered by Goldstone-limit unitarization

Summary

- ▶ Effective theory: good for TGC, limited applicability for QGC.
- ▶ Unitarization schemes tend to introduce theoretical prejudice
- ⇒ We propose a framework how to reconcile EFT with unitarity without losing its benefits
- ⇒ Direct T-Matrix unitarization as catch-all scheme for new models

Summary

- ▶ Effective theory: good for TGC, limited applicability for QGC.
- ▶ Unitarization schemes tend to introduce theoretical prejudice
- ⇒ We propose a framework how to reconcile EFT with unitarity without losing its benefits
- ⇒ Direct T-Matrix unitarization as catch-all scheme for new models

- ▶ Possible Realization: generic resonances = simplified model.

Summary

- ▶ Effective theory: good for TGC, limited applicability for QGC.
- ▶ Unitarization schemes tend to introduce theoretical prejudice
- ⇒ We propose a framework how to reconcile EFT with unitarity without losing its benefits
- ⇒ Direct T-Matrix unitarization as catch-all scheme for new models

- ▶ Possible Realization: generic resonances = simplified model.

- ▶ Extended Framework for quantitative tests of the SM version of electroweak interactions
- ⇒ Implemented in simulation for LHC/SPPC and ILC/CLIC in WHIZARD

Summary

- ▶ Effective theory: good for TGC, limited applicability for QGC.
- ▶ Unitarization schemes tend to introduce theoretical prejudice
- ⇒ We propose a framework how to reconcile EFT with unitarity without losing its benefits
- ⇒ Direct T-Matrix unitarization as catch-all scheme for new models

- ▶ Possible Realization: generic resonances = simplified model.

- ▶ Extended Framework for quantitative tests of the SM version of electroweak interactions
- ⇒ Implemented in simulation for LHC/SPPC and ILC/CLIC in WHIZARD