

INTRINSIC TIME QUANTUM GRAVITY & THE UNIQUE INITIAL STATE OF THE UNIVERSE

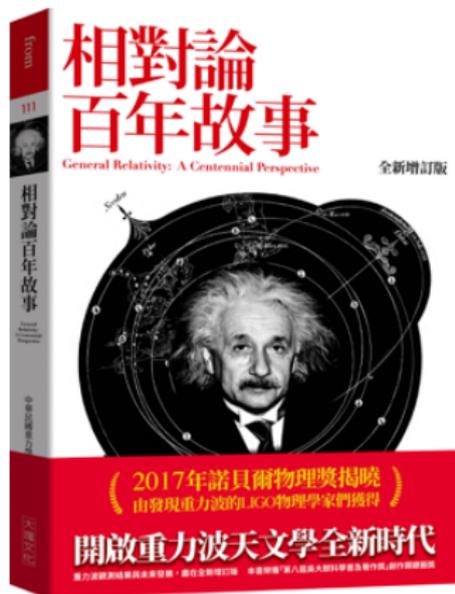
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GR: A Centennial Perspective:



The purposes of this talk:



What is the Most Important Observation about our Universe?

Universe's expansion should be taken fundamentally

- What are the majors fallacies of conventional inflation models ?

Outlines of this talk:

- 1 Conceptual and technical technical difficulties of GR in FWR Cosmology
 - Weyl tensor drops out from GR
- 2 Inflation doesn't solve the initial value problem of the Universe
 - Common features of inflationary models
- 3 Intrinsic Time Quantum Gravity(ITQG)
 - Singularity is necessary at the Big Bang to have Time
- 4 The initial state of the Universe is determined by the Cotton-York Tensor
- 5 ITQG ground state driven inflation without Inflaton & Experimental tests

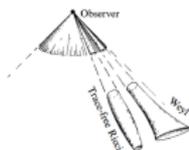
Universe Expansion should be noted fundamentally

- 1 The expansion of the Universe should be treated as fundamental as:
 - Newton's discovery of equivalence between falling apple and Moon's falling; therefore law of Gravitation
 - Einstein's discovery of equivalence between acceleration and gravitation; therefore GR
- 2 Universe's expansion means Time is special and therefore no space-time symmetry
- 3 Taking Universe's expansion seriously at fundamental level allows both a self-consistent formulation of QG and the existence of semi-classical limit

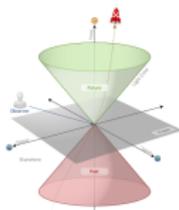
Major fallacies of GR & Inflationary models

- 1 Complete geometrical data of space-time needs the full 20-components Riemann Tensor;

$$R^{\mu}_{\nu\alpha\beta} = \text{Ricci}(\text{trace, source for gravity}) + \text{Weyl}(\text{traceless, dynamical d.o.f of gravity})$$



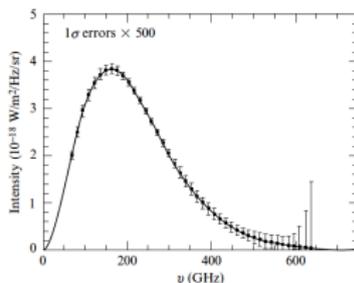
- 2 Eq. that det'd the Conformal Invariant Structure(Causal) (Weyl tensor) is missing from GR



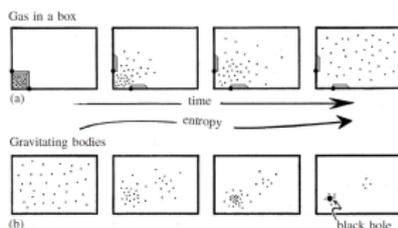
- 3 Matter can't tell spacetime how to curve; i.e. even in vacuum, spacetime form black holes
- 4 LIGO discovery of GW1050914 confirms that GW carries NON-ZERO local energy density
- 5 Homogeneity & isotropy conflicts the basic feature of gravitation, i.e. attractive
- 6 Wrong to describe inflation, a particular property of space-time by scalar matter fields
- 7 Predictions of inflationary models can be changed by varying the shape of the inflationary energy density curve or the initial conditions; **basically means no prediction**

Difference between Matter & Gravitation Entropy

- 1 Observation of 3K thermal(equilibrium) radiation \Rightarrow U. can't evolve after recombination!?



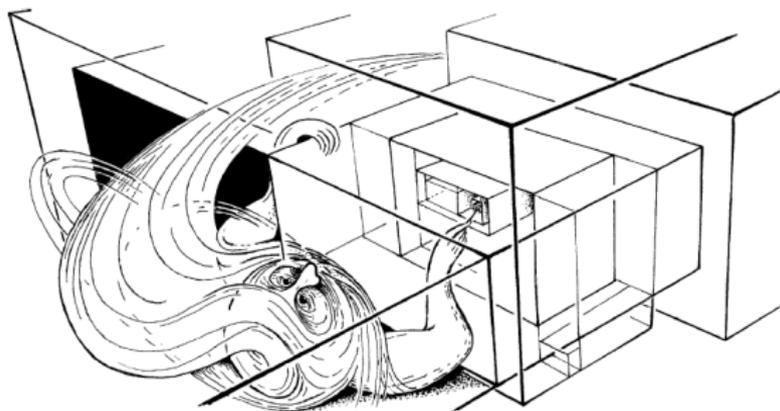
- 2 As matter homogenises, entropy increases but gravitation entropy increases when clumps



- 3 Gravity's d.o.fs are not in the thermal equilibrium(**means very special**) near Big Bang

The Initial state of the Universe is the Biggest puzzle

- 1 Entropy of solar mass BH is $\sim 10^{21}$, there about 10^{80} baryon in the Universe
- 2 Pro. of U. begin in such a low entropy state, $10^{-10^{123}}$; Cosmological const. is $\sim 10^{-120}$
- 3 Homogenous & isotropic (zero entropy) state in gravitation is very rare and special



Inflation is irrelevant in solving the initial state puzzle

- 1 The uniformity in the universe we now observed should be the result of inflation acting on its early evolution is basically misconceived
- 2 Homogenous and isotropic condition is still a MUST for inflation to happen; i.e. the basic Friedmann eqt. must be kept intacted!
- 3 Computer simulations requires homogenous within 3 Hubble lengths and vacuum dominated condition to have inflation to happen
- 4 There must be a great uniformity in the matter distribution at that time, for otherwise inflation cannot be introduced; A puzzle!?
- 5 Inflation only pushed the puzzle to an earlier epoch but didn't solve the uniformity puzzle
- 6 Inflation, a property of space-time structure evolution; shouldn't be driven by matter fields
- 7 Need to kill gravity entropy(not been thermalised) but keep its d.o.f. in the initial state

Conventional wisdoms about Physics of Inflation

- ① Starobinsky-like model captures the conformal structure through extra $\frac{R^2}{6M^2}$
- ② Going to Einstein frame through $q_{\mu\nu} \rightarrow e^{\phi} \tilde{g}_{\mu\nu}$; and obtain a characteristic potential with a large Mass scale dictated by M , i.e.

$$V = \tilde{R} + (\partial_{\mu}\phi)^2 - \frac{3}{2}M^2(1 - e^{\frac{\sqrt{2}}{3}\phi})$$

- ③ The appearance of this large Mass scale($M \gtrsim 10^{13} GeV$); drives exponential spatial expansion(inflation) in the Friedmann equation
- ④ Basically, all inflation models with scalar field inflaton MAKE use of this large mass scale
- ⑤ The problem is still requiring homogeneity to simplify the Einstein eqt. to Friedmann eqt.
- ⑥ All these models only push the initial condition to an earlier time; none of them explains: [why the homogenous and isotropy initial condition can be achieved at the Big Bang](#)

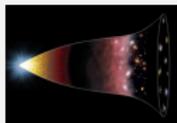
Diffeomorphism & Yang-Mills gauge structures



Table 1. Comparison of the two different gauge structures

	Diffeomorphism Gauge Structures	Yang-Mills Gauge Structures
Basic Variables	Spatial metric tensor q_{ij}	Gauge connection A_{ia}
Symmetry Generators	$H_i(\mathbf{x}) = -2q_{ik}\nabla_j\pi^{jk}(\mathbf{x})(=0)$	$G^a(\mathbf{x}) = \nabla_i\pi^{ia}(\mathbf{x})(=0)$
Gauge transformation	$[q_{ij}(\mathbf{x}), H_k[N^k]] = \mathcal{L}_{\vec{N}}q_{ij}(\mathbf{x});$ $H_i[N^i] = \int N^i H_i d^3\mathbf{x}$	$[A_{ia}(\mathbf{x}), G^b[\eta_b]] = -\nabla_i\eta_a(\mathbf{x})G^b[\eta_b];$ $G^b[\eta_b] = \int \eta_b G^b d^3\mathbf{x}$
Commutation Relations	$[H_i(\mathbf{x}), H_j(\mathbf{y})]$ $= H_j(\mathbf{x})\partial_i\delta(\mathbf{x}-\mathbf{y}) + H_i(\mathbf{y})\partial_j\delta(\mathbf{x}-\mathbf{y})$	$[G^a(\mathbf{x}), G^b(\mathbf{y})]$ $= if^ab_c G^c(\mathbf{x})\delta(\mathbf{x}-\mathbf{y})$
Potentials	$V \sim [\frac{\delta \exp(CS)}{\delta q_{ij}}]^2$	$V \sim [\frac{\delta \exp(CS)}{\delta A_{ia}}]^2$
Locality & Dimension	Not product of identical group(i.e. $SL(3R)$) at each spatial point of base manifold i.e. not of principle fibre structure	Infinite tensor product group $\prod_{\mathbf{x}} G;$ G =finite dimensional Lie group =usually referred to as the 'gauge group'

Initial Time Quantum Gravity



- ① Decompose $(q_{ij}, \tilde{\pi}^{ij})$ irreducibly into conjugate pairs: $(\bar{q}_{ij}, \bar{\pi}^{ij})$, $(\ln q^{\frac{1}{3}}, \pi)$; identifying $\ln q(x)$ as intrinsic time, $\pi = q_{ij} \tilde{\pi}^{ij}$ the energy function
- ② ADM Hamiltonian constraint ($\beta^2 = \frac{1}{6}$ for GR) delivers dynamical; not gauge evolution
 - $H(x) = -qR + \bar{q}_{ik} \bar{q}_{jl} \bar{\pi}^{ij} \bar{\pi}^{kl} - \beta^2 \pi^2 = 0$; doesn't generate Temporal gauge symmetries, otherwise, Universe becomes block diagonal & not consistent with QM
 - Dynamical eqt & local Hamiltonian density are derived from:

$$-\pi = \frac{\bar{H}(x)}{\beta} = \frac{1}{\beta} \sqrt{\bar{q}_{ik} \bar{q}_{jl} \bar{\pi}^{ij} \bar{\pi}^{kl} - qR} \rightarrow \frac{1}{\beta} \sqrt{\bar{q}_{ik} \bar{q}_{jl} \bar{\pi}^{ij} \bar{\pi}^{kl} + V}$$
 - Allow modifications of gravitation potential away from Einstein's theory
 - Classically, $H = K.E. + V$; Relativistically, $H = \sqrt{momenta^2 + V}$
 - The square root Hamiltonian correctly captures the dispersion relation $v = \frac{p}{E}$ & particle-wave duality with upper limit in speed; readily for carrying out quantization

Singularity is necessary at the Big Bang to have Time

- 1 What is diverging at the Big Bang is $\ln q = T \rightarrow -\infty$, which is exactly what we need to represent the beginning of Time in the Universe
- 2 Although R diverges (however, is suppressed by the Time factor) but \bar{q}_{ij} is non-zero by construction and the Cotton-York (C.Y.) tensor (density) is regular at the Big Bang
- 3 This necessary **temporal singularity** is usually misinterpreted as a curvature singularity
- 4 Gravitation d.o.fs are well defined and described by the C.Y. tensor at the Big Bang
- 5 The local Hamiltonian density stays semi-positive definite during the Big Bang
- 6 Big Bang singularity is necessary to allow the Universe to begin

Conformal flatness initial state & Penrose Conjecture

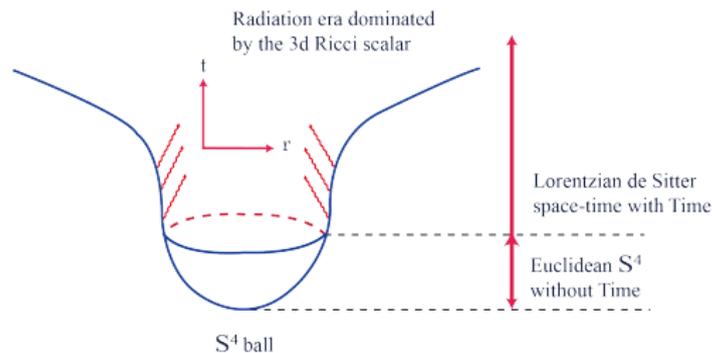
- 1 $\bar{H} = \sqrt{(\text{extrinsic curvature})^2 + (C.Y. tensor)^2 + q\mathcal{K}} = \sqrt{\bar{\pi}_j^i \bar{\pi}_i^j + (C.Y. tensor)^2 + q\mathcal{K}}$
- 2 \bar{H} attains classical minimum iff $\bar{\pi}_j^i$ & C.Y. vanishes identically, precisely the criterion for conformal flatness(S^3) in 3d
- 3 Vanishing of the extrinsic curvature components & spatial conformal flatness at $T \rightarrow -\infty$ together with $\nabla_i \bar{\pi}^{ij} = 0$, realizes a FRW Big Bang compatible with 4d Weyl tensor = 0
- 4 'Our extraordinarily special Big Bang' with low entropy emerges naturally from the ground state of $\bar{H} = 0$ in the C.Y. tensor dominance era & rules out the possibilities of primordial black hole
- 5 2nd Law thermodynamic 'arrow of time' & 'gravitational arrow of intrinsic time'(increasing vol.) point in same direction
- 6 Nontrivial realisation of Poincare conjecture; ∞ nos. 3d simply connected diff. structures all conformally equiv. to S^3

The initial Quantum State

- 1 C.Y. dominating era(right after Big Bang), $\bar{H} = \sqrt{\hat{\pi}_j^{\dagger i} \hat{\pi}_i^j + g^2 \hbar^2 \tilde{C}_j^i \tilde{C}_i^j + q\mathcal{K}}$; $[\hat{\pi}_i^j, \tilde{C}_i^j] = 0$
- 2 $\hat{\pi}_j^i \rightarrow \hat{Q}_j^i = e^{gW_T} \hat{\pi}_j^i e^{-gW_T} = \frac{\hbar}{i} E_j^{i(mn)} [\frac{\delta}{\delta \bar{q}_{ij}} - g \frac{\delta W_T}{\delta \bar{q}_{mn}}] = \hat{\pi}_j^i + ig\hbar \tilde{C}_j^i$; $W_T = W_{CS} + W_{EH}$
 $= \frac{g}{4} \int \epsilon^{ijkl} (\Gamma_{im}^l \partial_j \Gamma_{kl}^m + \frac{2}{3} \Gamma_{im}^l \Gamma_{jn}^m \Gamma_{kl}^n) d^3x - \alpha \int \sqrt{q} R d^3x$; $Q_j^i, Q_i^{\dagger j}$ are unitarily inequivalent SL(3R) reps.
- 3 $\tilde{C}_j^i = \frac{\delta W_{CS}}{\delta q_{ij}}$; is the C.Y. tensor(density); analogous to the magnetic field in the Yang-Mills Hamiltonian density,
 $\frac{q_{ij}}{2} (E^{ia} E_a^j + B^{ia} B_a^j) = \frac{q_{ij}}{2} \hat{Q}^{\dagger ia} \hat{Q}^j a$; $\hat{Q}^j a = e^{W_{CS}} E^{ia} e^{-W_{CS}} = E^{ia} + iB^{ia}$ & $B^{ia} = \frac{\delta W_{CS}}{\delta A_{ia}}$
- 4 Hamiltonian density becomes $\bar{H} = \sqrt{\hat{Q}_j^{\dagger i} \hat{Q}_i^j + q\mathcal{K}}$ with $\hat{Q}_j^i e^{\pm gW_T} |0\rangle = 0$; $\hat{Q}_j^{\dagger i}$ and \hat{Q}_i^j are not symmetrical;
 the generalised coherent state generated is given by $e^{\int \frac{1}{\sqrt{q}} \hat{C} \cdot \hat{Q}^{\dagger}} |0\rangle$
- 5 Classical conformally flat(C.Y.=0) is an extremum of W , and thus precisely a saddle point for the quantum ground state wavefunctional $\Psi_0[\bar{q}_{ij}] = N e^{g(W - W_o)} \langle \bar{q}_{ij} | 0 \rangle$
- 6 Imaginary momentric means $\Psi_0[\bar{q}_{ij}]$ is a quantum tunnelling solution, however, for compact manifolds, i.e. S^3 ;
 $\langle \Psi_0 | \hat{\pi}_j^{\dagger i} | \Psi_0 \rangle = 0$, $\langle \Psi_0 | \tilde{C}_j^i | \Psi_0 \rangle = 0$ (metric can still fluctuate; steepest descent)
 $\Rightarrow \Psi_0[\bar{q}_{ij}]$ is barely classically allowed; therefore the initial state indeed satisfies the Weyl hypothesis
- 7 Important feature: Q Time=Semiclassical Time=Classical Time; Asymptotic free quantum state is also a semiclassical state \Rightarrow classical world is always possible
- 8 Intrinsic time framework allows continuation of β across zero from real to imaginary values(Lorentzian to Euclidean)

Inflation by ground state property without Inflaton

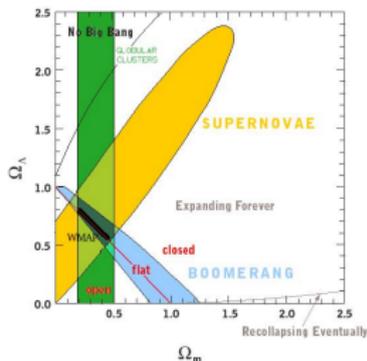
- 1 $\langle \hat{\pi}_j^i \rangle = 0$, force metric fluctuations buried inside $\langle \tilde{C}_j^i \rangle = 0 \Rightarrow \langle \pi \rangle \propto \langle \bar{H} \rangle = 0$ (allow $\beta \rightarrow 0$)
 \Rightarrow zero extrinsic curvature, the junction condition needed for Euclidean to Lorentzian continuation of the metric
- 2 Exponential growth in spatial volume as time changes is a result forced by geometries of the initial state at the throat of the de Sitter junction after the Big Bang. No need for scalar field inflaton!



The equator of S^4 is S^3 . The de Sitter manifold joins to this S^3 ball tangentially because of the vanishing extrinsic curvature. Therefore the spatial exponential expansion is due to the geometric junction conditions.

Observational Tests of the New Model

- 1 Zero extrinsic curvature & C.Y. tensor exclude the possibility of $k = 0$, spatially flat \mathcal{R}^3 & $k = -1$, spatially open \mathcal{H}^3
- 2 This will be confirmed by the future high precision mass density measurements



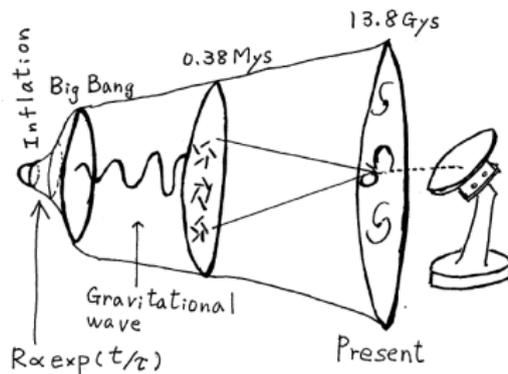
- 3 The de Sitter expansion period is controlled by $T_f - T_i = \frac{2}{3} \ln \frac{V_f}{V_i}$; $a_f = a_i e^N$ with $N = \frac{1}{2}(T_f - T_i) \sim 50 - 60$ e-folding before the Ricci scalar terms begin to dominate
- 4 Neither scalar field inflaton nor large v.e.v. is needed to drive the Universe into the de Sitter expansion period, instead, the self-determined geometries fix all the details of the Universe's evolution/expansion at the beginning of the Big Bang

Predictions of primordial scale invariant tensor modes

- 1 C.Y.=0 is the classical extremum of W , & thus precisely a saddle point for the quantum ground state wavefunction; $\Psi_0[\bar{q}_{ij}] = N e^{g(W-W_0)} \langle \bar{q}_{ij} | 0 \rangle$
- 2 The Quantum ground state therefore contains primordial fluctuations (CS is a topology inv. entity under infinitesimal variation of metric); expansion of W about the saddle point $\frac{\delta W}{\delta \bar{q}_{ij}} = \tilde{C}^{ij} = 0 \Rightarrow$ primordial tensor mode fluctuations

$$W_{\bar{q}_{ij}} - W_0 = W_{\tilde{C}^i_j=0} + \frac{1}{2} \int d^3x \int d^3y \delta \bar{q}_{ij}(x) H^{ijkl}(x, y) |_{\tilde{C}^m_n=0} \delta \bar{q}_{kl}(y) + \dots - W_0$$

- 3 Predicts $\sim 1/k^3$ scale invariance tensor mode fluctuations; can be measured by the BICEP experiments



Quantum Hamiltonian

$$\textcircled{1} \quad \bar{H} = \sqrt{Q_j^{\dagger i} Q_i^j + q\mathcal{K}} = \sqrt{\bar{\pi}_i^{\dagger j} \bar{\pi}_j^i + \hbar^2 (g\tilde{C}_j^i - \alpha\sqrt{q}\bar{R}_j^i)(g\tilde{C}_i^j - \alpha\sqrt{q}\bar{R}_i^j) - i\alpha\hbar\sqrt{q}[\bar{\pi}_j^i, \bar{R}_i^j] + q\mathcal{K}};$$

$$\tilde{C}_j^i = \bar{E}_{j(mn)}^i \frac{\delta W_{CS}}{\delta \bar{q}_{mn}}, \quad W_{CS} = \frac{g}{4} \int \tilde{\epsilon}^{ijk} (\Gamma_{im}^l \partial_j \Gamma_{kl}^m + \frac{2}{3} \Gamma_{im}^l \Gamma_{jn}^m \Gamma_{kl}^n) d^3x \quad \& \quad [\bar{\pi}_j^i, \tilde{C}_i^j] = 0$$

- $\textcircled{2}$ $Q_j^{\dagger i}$ & Q_i^j related to $\bar{\pi}_j^i$ by $e^{\mp W_T}$, are non-Hermitian, generate 2 unitarily inequivalent rep. of the non-compact group $SL(3, R)$ at each x ; $\bar{\pi}_j^i = \frac{1}{2}(Q_j^{\dagger i} + Q_i^j)$ generates a unitary rep. of $\prod_x SU(3)_x$. In fact, the commutator term in \bar{H} can be identified as $-i\alpha\hbar\sqrt{q}[\bar{\pi}_j^i, \bar{R}_i^j] = [Q_j^{\dagger i}, Q_i^j]$

- $\textcircled{3}$ The commutator at coincident spatial position in functional Schrodinger representation is,

$$[Q_j^{\dagger i}(x), Q_i^j(x)] = -\alpha\hbar^2 \sqrt{q} \bar{E}_{jkl}^i(x) \frac{\delta \bar{R}_i^j(x)}{\delta \bar{q}_{kl}(x)} = -\alpha\hbar^2 \sqrt{q(x)} \int d^3y \delta(x-y) E_{jkl}^i(x) \frac{\delta R_i^j(y)}{\delta q_{kl}(x)}$$

- $\textcircled{4}$ Variation of the Ricci tensor;

$$\begin{aligned} \delta R_{mn} &= -\frac{1}{2}(\nabla^2 + R)\delta q_{mn} + \frac{3}{2}(R_n^s \delta q_{sm} + R_m^s \delta q_{sn}) + \frac{1}{2}(\nabla_m \nabla^s \delta q_{sn} + \nabla_n \nabla^s \delta q_{sm}) \\ &\quad - (R_{mn} - \frac{1}{2}q_{mn}R + \frac{1}{2}\nabla_m \nabla_n) \delta \ln q \end{aligned}$$

Emergence of Einstein's Theory I

- 1 Through integration by parts, the point split commutator term becomes

$$[Q_j^\dagger(x), Q_i^j(y)] = -\frac{5}{6}\alpha\hbar^2\sqrt{q(x)}(-\nabla_x^2 + R)\delta(x-y). \quad (1)$$

- 2 Heat kernel, $K(\epsilon; x, y)$ with $\lim_{\epsilon \rightarrow 0} K(\epsilon; x, y) = \delta(x-y)$, regulates the coincident limit in (1) for generic metrics;

$$\nabla^2 K(\epsilon; x, y) = \frac{\partial K(\epsilon; x, y)}{\partial \epsilon};$$

Seeley-DeWitt coefficients a_n , $2\sigma(x, y)$ square of the geodesic length, & Δ_V the Van Vleck determinant,

$$K(\epsilon; x, y) = (4\pi\epsilon)^{-\frac{3}{2}}\Delta_V^{\frac{1}{2}}(x, y)e^{-\frac{\sigma(x, y)}{2}\epsilon}\sqrt{q(y)}\sum_{n=0}^{\infty}a_n(x, y; \nabla^2)\epsilon^n,$$

wherein ϵ is of dimension L^2 . In the coincidence $x = y$ limit, the coefficients for closed manifolds are

$$a_0 = 1; \quad a_1 = \frac{R}{6}; \quad a_2 = \frac{1}{180}(R^{ijkl}R_{ijkl} - R^{ij}R_{ij}) + \frac{1}{30}\nabla^2 R + \frac{R^2}{72}$$

- 3 Coincidence limit of $\lim_{\epsilon \rightarrow 0} K(\epsilon; x, y) = \delta(x-y)$ regulates $\delta(0)$, and $\lim_{\epsilon \rightarrow 0} \frac{\partial K(\epsilon; x, x)}{\partial \epsilon}$ for $-\nabla^2 \delta(0)$,

Emergence of Einstein's Theory II

$$\textcircled{1} [Q_j^{\dagger i}(x), Q_i^j(x)] = \lim_{\epsilon \rightarrow 0} -\frac{5\alpha\hbar^2 q}{6(4\pi\epsilon)^{\frac{3}{2}}} \left[\frac{3}{2\epsilon} + \frac{13R}{12} + (a_1 R - \frac{a_2}{2})\epsilon + \text{terms with } \epsilon^{n \geq 2} \right];$$

$\textcircled{2}$ This yields the full fledged local Hamiltonian density

$$\bar{H} = \lim_{\epsilon \rightarrow 0} \sqrt{\pi_i^{\dagger j} \pi_j^i + \hbar^2 (g\tilde{C}_j^i - \alpha\sqrt{q}\tilde{R}_j^i)(g\tilde{C}_i^j - \alpha\sqrt{q}\tilde{R}_i^j) + q\left(\mathcal{K} - \frac{5\alpha\hbar^2}{32\pi^{\frac{3}{2}}\epsilon^{\frac{5}{2}}}\right) - \frac{5\alpha\hbar^2 q}{6(4\pi\epsilon)^{\frac{3}{2}}} \left(\frac{13}{12}R + (a_1 R - \frac{a_2}{2})\epsilon\right)};$$

3 fundamental coupling constants g , \mathcal{K} and α . Upon regulator removal, $\epsilon \rightarrow 0$, divergences are countered by \mathcal{K} and α

$\textcircled{3}$ Renormalized values are identified phenomenologically, with $\frac{1}{(2\kappa)^2} = \frac{5\alpha\hbar^2}{6(4\pi\epsilon)^{\frac{3}{2}}} \left(\frac{13}{12}\right)$ to yield the correct Newtonian limit, and read off the effective cosmological constant Λ_{eff} from $\frac{2\Lambda_{\text{eff}}}{(2\kappa)^2} = \left(\mathcal{K} - \frac{5\alpha\hbar^2}{32\pi^{\frac{3}{2}}\epsilon^{\frac{5}{2}}}\right) = \left(\mathcal{K} - \frac{2^4 3^2}{13\kappa^2 \epsilon}\right)$

$\textcircled{4}$ $\alpha = \frac{72(4\pi\epsilon)^{\frac{3}{2}}}{65\hbar^2(2\kappa)^2}$, and finiteness of κ implies $\alpha \rightarrow 0$ as $\epsilon \rightarrow 0$. The theory actually produces all the Seeley-DeWitt

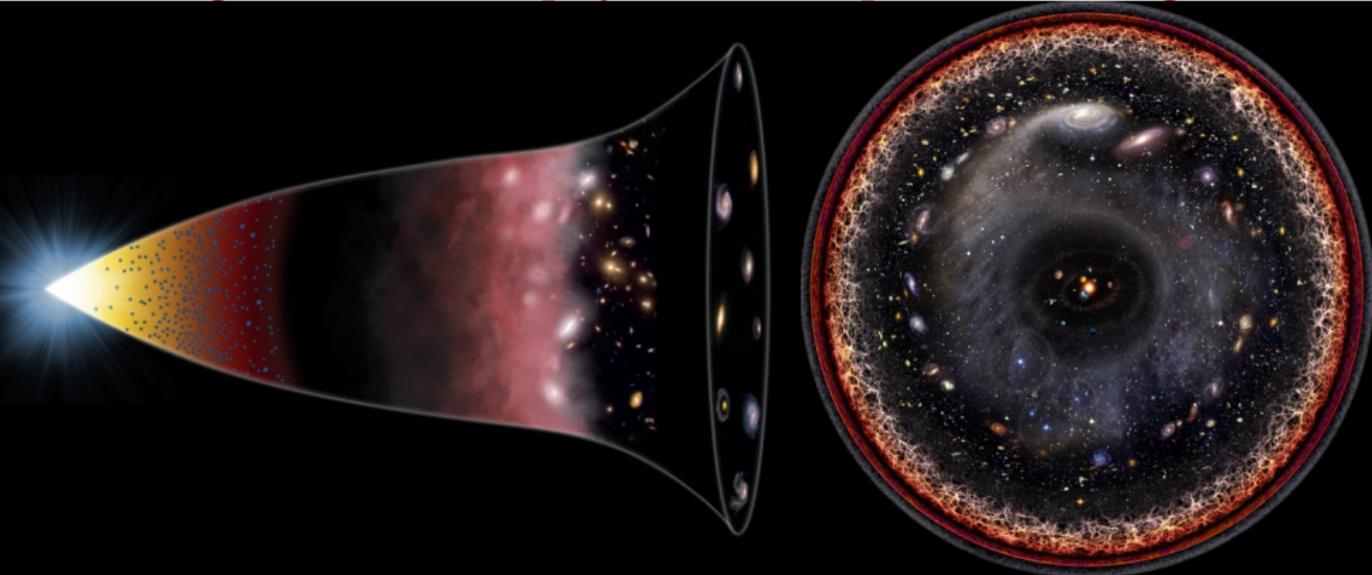
coefficients, but these higher curvature terms with positive powers of ϵ all disappear upon regulator removal

$\textcircled{5}$ Finally $\bar{H} = \sqrt{Q_j^{\dagger i} Q_i^j + q\mathcal{K}} = \sqrt{\pi_i^{\dagger j} \pi_j^i + \hbar^2 g^2 \tilde{C}_j^i \tilde{C}_i^j - \frac{q}{(2\kappa)^2} (R - 2\Lambda_{\text{eff}})}$

Four Pillars of whys, that make the Universe as it is

- ① **Causality:** Foundation of orders \rightarrow allows our mother nature being accessible to human minds
 - Gravity forces $\{, \}_{P.B.}$ to be replaced by $[,] \rightarrow$ Causality regained dynamically
 - Realised by a single Ψ (therefore can globally be agreed) in Quantum Mechanics
- ② **Unitary:** Information preserving, none will be left behind \rightarrow Universe to evolve as a whole inside the Ψ
- ③ **Computability:** Dimensionless Schrödinger eqt. (\hbar & κ being conversion factors) unifies everything into a single GI Ψ
 - $\{, \}_{P.B.}$ replaced by $[,] \rightarrow$ evolution guaranteed; Universe needs no reboot
 - Gauged \rightarrow To get rid of the ∞ d.o.fs in the field operator for 'Time' \rightarrow integrability
 - Completeness \rightarrow 1^{st} order S-eqt. guarantees consistency with trillions of H-J eqt.
 - Integrability & Complexity \rightarrow only **ONE** Time & flows in the direction of entropy \uparrow
- ④ **Renormalizability:** Allow structureless point-like Quantum Fields to make Geometroynamics possible
 - Each point between the Small/Large Universe can be identified
 - Geometroynamics allows the Universe to expand without leaving out anything

Meaning of Time - A physicist's response to Bergson



Time and Causality is to evoke existence: Things come into beings and facts become true all the time & therefore flows in the same direction as entropy↑

Meaning requests causality & the reality of time, therefore must be QM